## ME/CS 133(a): Solution to Homework \#1

(Fall 2018/2019)

## Solution to Problem 1:

Let the $2 \times 1$ vectors ${ }^{1} \vec{v}=\left[\begin{array}{ll}{ }^{1} v_{1} & { }^{1} v_{2}\end{array}\right]^{T}$ and ${ }^{2} \vec{v}=\left[\begin{array}{ll}{ }^{2} v_{1} & { }^{2} v_{2}\end{array}\right]^{T}$ have associated complex representations ${ }^{1} \tilde{v}={ }^{1} v_{1}+i^{1} v_{2}$ and ${ }^{2} \tilde{v}={ }^{2} v_{1}+i^{2} v_{2}$ respectively (where $i^{2}=-1$ ). Recall that the goal of this problem is to show that the complex number formula:

$$
\begin{equation*}
{ }^{1} \tilde{v}=\tilde{d}_{12}+e^{i \theta_{12}}{ }^{2} \tilde{v} \tag{1}
\end{equation*}
$$

is equivalent to the planar coordinate transformation:

$$
\begin{equation*}
{ }^{1} \vec{v}=\vec{d}_{12}+R\left(\theta_{12}\right)^{2} \vec{v} . \tag{2}
\end{equation*}
$$

Let's evaluate the right hand side of expression (1) using the standard rules for multiplication of complex numbers ${ }^{1}$ :

$$
\begin{align*}
\tilde{d}_{12}+e^{i \theta_{12}}{ }^{2} \tilde{v} & =(x+i y)+\left(\cos \theta_{12}+i \sin \theta_{12}\right)\left({ }^{2} v_{1}+i^{2} v_{2}\right) \\
& =\left(x+{ }^{2} v_{1} \cos \theta_{12}-{ }^{2} v_{2} \sin \theta_{12}\right)+i\left(y+{ }^{2} v_{1} \sin \theta_{12}+{ }^{2} v_{2} \cos \theta_{12}\right) \tag{3}
\end{align*}
$$

where we have used Euler's formula $\left(e^{i \theta}=\cos \theta+i \sin \theta\right)$. Matching the real and complex portions of Equation (3) with the real and complex parts of ${ }^{1} \tilde{v}$ in the left hand side of Equation (1), we see that

$$
\begin{align*}
& { }^{1} v_{1}=x+{ }^{2} v_{1} \cos \theta-{ }^{2} v_{2} \sin \theta  \tag{4}\\
& { }^{1} v_{2}=y+{ }^{2} v_{1} \sin \theta+{ }^{2} v_{2} \cos \theta \tag{5}
\end{align*}
$$

These equations are equivalent to

$$
{ }^{1} \vec{v}=\vec{d}_{12}+\left[\begin{array}{cc}
\cos \theta_{12} & -\sin \theta_{12}  \tag{6}\\
\sin \theta_{12} & \cos \theta_{12}
\end{array}\right]{ }^{2} \vec{v}
$$

Solution to Problem 2: Recall that the location of the pole is fixed in both the moving and observer reference frames. Hence, before displacement, the pole is located at some position ${ }^{B} \vec{p}$ as seen by an observer in the fixed $B$ frame. After displacement, the observer in the body fixed $C$ frame also sees the pole in his/her coordinates at point ${ }^{B} \vec{p}$. However, the moving body has displaced relative to the fixed observer by amount $D_{12}=\left(\vec{d}_{12}, R_{12}\right)$. But points in the observer and displaced reference frames are related by a coodinate transform. Since the pole is at the same location in both the fixed and moving frames, it must be true that:

$$
{ }^{B} \vec{p}=\vec{d}_{12}+R_{12} \quad{ }^{B} \vec{p} .
$$

This equation can be solved to find the pole location:

$$
{ }^{B} \vec{p}=\left(I-R_{12}\right)^{-1} \vec{d}_{12}
$$

[^0]Of course, the matrix ( $I-R_{12}$ ) must be invertible, which will alwas be true except when $R_{12}=I$. In this case, the motion is a pure translation, which is viewed as a rotation about the "pole at infinity."
B) In Frame B, the pole is located at: ${ }^{B} \vec{p}=\left(I-R_{12}\right)^{-1} \vec{d}_{12}$
C) In Frame $C$, the vector describing the pole has exactly the same value as seen by the observer in Frame B: ${ }^{C} \vec{p}=\left(I-R_{12}\right)^{-1} \vec{d}_{12}$
A) In Frame A, the expression for the pole vector is obtained by a simple coordinate transformation of the expression in Frame B: ${ }^{A} \vec{p}=\vec{d}_{01}+R_{01}{ }^{B} \vec{p}=\vec{d}_{01}+R_{01}\left(I-R_{12}\right)^{-1} \vec{d}_{12}$

Problem 3: To find the pole of the displacement: $D_{2}=(x, y, \theta)=\left(3.0,2.0,45.0^{\circ}\right)$, substitute into the above results:

$$
\begin{align*}
{ }^{B} \vec{p}=\left(I-R_{12}\right)^{-1} \vec{d}_{12} & =\left[\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right)-\left(\begin{array}{cc}
\cos \left(45^{\circ}\right) & -\sin \left(45^{\circ}\right) \\
\sin \left(45^{\circ}\right) & \cos \left(45^{\circ}\right)
\end{array}\right)\right]^{-1}\left[\begin{array}{l}
3.0 \\
2.0
\end{array}\right] \\
& =\left[\begin{array}{cc}
1-\frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \\
-\frac{\sqrt{2}}{2} & 1-\frac{\sqrt{2}}{2}
\end{array}\right]^{-1}\left[\begin{array}{l}
3.0 \\
2.0
\end{array}\right] \\
& =\left[\begin{array}{l}
? \\
?
\end{array}\right] \tag{7}
\end{align*}
$$

You could report this result in Frame B, or transform the results to frame A.

$$
\begin{align*}
{ }^{A} \vec{p}=\vec{d}_{01}+R_{01} \quad{ }^{B} \vec{p} & =\left[\begin{array}{l}
2.0 \\
2.0
\end{array}\right]+\left(\begin{array}{cc}
\cos \left(20^{\circ}\right) & -\sin \left(20^{\circ}\right) \\
\sin \left(20^{\circ}\right) & \cos \left(20^{\circ}\right)
\end{array}\right)\left[\begin{array}{l}
? \\
?
\end{array}\right]  \tag{8}\\
& =\left[\begin{array}{l}
? \\
?
\end{array}\right] \tag{9}
\end{align*}
$$

Problem 4: To show that a transformation is a pure rotation when viewed in a reference frame at the pole, select a new reference frame, denoted by $D$, whose basis vectors are parallel to Frame B and whose origin lies at the pole of the displacement. Let $\vec{p}$ denote the location of the pole, as seen by an observer in Frame B. The location of Frame $B$ relative to Frame $D$ is a pure translation of amount $-{ }^{1} \vec{p}$, and therefore, $D_{D B}=(-\vec{p}, I)$. The displacement of the body from the first position to the second position, as now observed in Frame $D$, is obtained by a similarity transform $D_{D B} D_{12} D_{D B}^{-1}$ :

$$
\begin{align*}
D_{D B} D_{12} D_{D B}^{-1} & =(-\vec{p}, I)\left(\vec{d}_{12}, R_{12}\right)(-\vec{p}, I)^{-1}  \tag{10}\\
& =(-\vec{p}, I)\left(\vec{d}_{12}, R_{12}\right)(+\vec{p}, I)  \tag{11}\\
& =(-\vec{p}, I)\left(\left(\vec{d}_{12}+R_{12} \vec{p}\right), R_{12}\right)  \tag{12}\\
& =\left(\left(\vec{d}_{12}+\left(R_{12}-I\right) \vec{p}\right), R_{12}\right) \tag{13}
\end{align*}
$$

Hence, if $\vec{p}=-\left(R_{12}-I\right)^{-1} \vec{d}_{12}=\left(I-R_{12}\right)^{-1} \vec{d}_{12}$, then $D_{D B} D_{12} D_{D B}^{-1}=\left(\overrightarrow{0}, R_{12}\right)$. I.e., as viewed in reference Frame $D$, the displacement is a pure rotation by amount $R_{12}$.

## Problem 5:

Part (a): There are many ways that one can prove that reflections preserve length. Here is one approach (see Figure 1).


Figure 1: Geometry of Planar Rigid Body Reflection
Select any two non-identical points, $A$ and $B$, in a rigid body. After reflection, those points become $A^{\prime}$ and $B^{\prime}$. Form the right triangle $A B D$, where the line $B D$ is chosen to be perpendicular to the line $A A^{\prime}$. Similary, in the reflected body, form the right triangle $A^{\prime} B^{\prime} D^{\prime}$. Simple geometric arguments show that since the distance $|B D|$ and $\left|B^{\prime} D^{\prime}\right|$ are equal, and the distances $|A D|$ and $\left|A^{\prime} D^{\prime}\right|$ are equal, then $|A B|=\left|A^{\prime} B^{\prime}\right|$. Hence, the distance between $A$ and $B$ is preserved under reflection. Since $A$ and $B$ were chosen randomly, the result will hold for any non-identical pair of points in the body. Thus, distance is always preserved under reflection.

Part (b): Generally, physically meaningful planar displacements are not equivalent to a single reflection. To see this, define three points $(A, B, C)$ in the body of Figure 1. Because the body is rigid, one can think of points $(A, B, C)$ as forming a rigid triangle. Consider the triangle formed from the reflected points $\left(A^{\prime}, B^{\prime}, C^{\prime}\right)$. Note that it is impossible physically translate $(A, B, C)$ to ( $A^{\prime}, B^{\prime}, C^{\prime}$ ). Finally, note that any rigid body planar displacement can generally be realized as the result of two sequential reflections.

Problem 6: (Problem 2.10(a,b,c) in MLS)
Part (b): If $\omega \in \mathbb{R}$, then let:

$$
\hat{\omega}=\left[\begin{array}{cc}
0 & -\omega \\
\omega & 0
\end{array}\right]=\omega J \quad \text { where } J=\left[\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right] .
$$

Note that:

$$
\hat{\omega}^{2}=-\omega^{2} I ; \quad \hat{\omega}^{3}=\omega^{3} J
$$

Hence:

$$
\begin{aligned}
e^{\hat{\omega} \theta} & =I+\omega \theta J+\frac{(\omega \theta)^{2}}{2!} J^{2}+\frac{(\omega \theta)^{3}}{3!} J^{3}+\cdots \\
& =I+(\omega \theta) J-\frac{(\omega \theta)^{2}}{2!} I-\frac{(\omega \theta)^{3}}{3!} J+\cdots \\
& =\left(1+\frac{(\omega \theta)^{2}}{2!}+\frac{(\omega \theta)^{4}}{4!}+\cdots\right) I+\left(\omega \theta-\frac{(\omega \theta)^{3}}{3!}+\cdots\right) J \\
& =\cos (\omega \theta) I+\sin (\omega \theta) J \\
& =\left[\begin{array}{cc}
\cos (\omega \theta) & -\sin (\omega \theta) \\
\sin (\omega \theta) & \cos (\omega \theta)
\end{array}\right]
\end{aligned}
$$

We can think of the map exp :so(2) $\rightarrow S O(3)$ as embedding a planar rotation into $S O(3)$, where we have chosen the axis of rotation to be perpendicular to the $x-y$ plane:

$$
\left[\begin{array}{ccc}
\cos (\omega \theta) & -\sin (\omega \theta) & 0 \\
\sin (\omega \theta) & \cos (\omega \theta) & 0 \\
0 & 0 & 1
\end{array}\right] .
$$

Clearly this map can not be surjective, since there are many rotations in $S O(3)$ which can not be expressed as the exponential of an element in so(2). If we assume that the product $\omega \theta$ is nonzero except at $\theta=0$ (i.e., assume that $\omega \neq 0$ and $\theta \in[-\pi, \pi]$ ), then the map is not injective since $\exp (\omega \theta)=\exp (\omega \theta+2 \pi)$.

Part (c): This can be shown by brute-force calculation. If $R \in S O(2)$ then:

$$
\begin{aligned}
R \hat{\omega} R^{T} & =\left[\begin{array}{cc}
\cos \theta & -\sin \theta \\
\sin \theta & \cos \theta
\end{array}\right]\left[\begin{array}{cc}
0 & -\omega \\
\omega & 0
\end{array}\right]\left[\begin{array}{cc}
\cos \theta & \sin \theta \\
-\sin \theta & \cos \theta
\end{array}\right] \\
& =\omega\left[\begin{array}{cc}
\cos \theta & -\sin \theta \\
\sin \theta & \cos \theta
\end{array}\right]\left[\begin{array}{cc}
\sin \theta & -\cos \theta \\
\cos \theta & \sin \theta
\end{array}\right] \\
& =\omega\left[\begin{array}{cc}
0 & -\left(\cos ^{2} \theta+\sin ^{2} \theta\right) \\
\left(\sin ^{2} \theta+\cos ^{2} \theta\right) & 0
\end{array}\right]=\omega\left[\begin{array}{cc}
0 & -1 \\
1 & 0
\end{array}\right]=\hat{\omega}
\end{aligned}
$$


[^0]:    ${ }^{1}$ If $\tilde{a}=a_{1}+i a_{2}$ and $\tilde{b}=b_{1}+i b_{2}$, then $\tilde{a} \tilde{b}=\left(a_{1} b_{2}-a_{2} b_{2}\right)+i\left(a_{1} b_{2}+a_{2} b_{1}\right)$.

